

EXAM BLOCK 2 GEOPHYSICAL FLUID DYNAMICS

17 March 2010, 9.00 - 11.00 hours

Two problems (all items have equal weight)

Remark 1: answers may be written in English or Dutch.

Remark 2: in all questions you may use $g = 10 \text{ ms}^{-2}$, $a = 6400 \text{ km}$ and $\Omega = 7.3 \times 10^{-5} \text{ s}^{-1}$.

Problem 1

A fluid system is governed by the following equations:

$$\begin{aligned}\frac{\partial u}{\partial t} - \beta_0 y v &= g' \frac{\partial a}{\partial x}, \\ \frac{\partial v}{\partial t} + \beta_0 y u &= g' \frac{\partial a}{\partial y}, \\ -\frac{\partial a}{\partial t} + H \left[\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right] &= 0.\end{aligned}$$

a. Name the set of these equations.

Also, describe the meaning of the symbols β_0 , g' , a and H .

b. When focusing on flow that evolves on timescales of many days, the system above can be reduced to a single equation for variable v , i.e.,

$$\frac{\partial}{\partial t} \left(\nabla^2 v - \frac{\beta_0^2 y^2}{g' H} v \right) + \beta_0 \frac{\partial v}{\partial x} = 0,$$

which admits wave-like solutions of the form

$$v = \Re \left\{ V(y) e^{i(kx - \omega t)} \right\}.$$

Derive the equation for $V(y)$.

c. The equation for $V(y)$ found in item b has a solution

$$V = V_0 y \exp \left(-\frac{1}{2} \mu^2 y^2 \right), \quad \mu^2 = \frac{\beta_0}{\sqrt{g' H}},$$

with V_0 a constant and the corresponding dispersion relation

$$\omega = \frac{-\beta_0 k}{k^2 + 3\mu^2}.$$

What kind of waves are described by this solution?

Also, give a physical interpretation of parameter μ .

d. Compute the range of wavenumbers k for which the phase velocity and group velocity of these waves are of opposite sign.

For problem 2: P.T.O.

Problem 2

Consider flow governed by

$$\frac{\partial q}{\partial t} + J(\psi, q) = 0,$$
$$q = \nabla^2 \psi + \frac{\partial}{\partial z} \left(\frac{f_0^2}{N^2} \frac{\partial \psi}{\partial z} \right) + \beta_0 y.$$

- What principle is expressed by this equation?
- Give a physical interpretation of the three terms that constitute variable q .
- Show that the equation above admits a solution that represents a steady, zonal flow with an arbitrary vertical structure.
- Assume a steady zonal flow in the atmosphere, which is governed by the equation given above, and which has a constant vertical shear $\alpha = 10^{-3} \text{ s}^{-1}$. Compute the magnitude and direction of the thermal wind between vertical levels $z = 1 \text{ km}$ and $z = 5 \text{ km}$.
- Compute the density field that corresponds to the flow considered in item d.

END

GFD 2009 Equation sheet

Continuity and momentum equations: molecular viscous fluid

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x}(\rho u) + \frac{\partial}{\partial y}(\rho v) + \frac{\partial}{\partial z}(\rho w) = 0$$

$$\begin{aligned}\rho \left(\frac{du}{dt} + f_* w - f v \right) &= - \frac{\partial p}{\partial x} + \mu \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \right) \\ \rho \left(\frac{dv}{dt} + f u \right) &= - \frac{\partial p}{\partial y} + \mu \left(\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} + \frac{\partial^2 v}{\partial z^2} \right) \\ \rho \left(\frac{dw}{dt} - f_* u \right) &= - \frac{\partial p}{\partial z} - \rho g + \mu \left(\frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial y^2} + \frac{\partial^2 w}{\partial z^2} \right)\end{aligned}$$

Energy budget for adiabatic flow of fixed composition

$$\rho C_v \frac{dT}{dt} - \frac{T}{\rho} \left(\frac{\partial p}{\partial T} \right)_\rho \frac{d\rho}{dt} = 0$$

Relative circulation and relative vorticity

$$\Gamma = \oint_C \mathbf{u} \cdot d\mathbf{r} = \int_S (\nabla \times \mathbf{u}) \cdot \mathbf{n} dS$$

where S is the surface enclosed by contour C .

Shallow water equations

$$\begin{aligned}\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} - f v &= -g \frac{\partial \eta}{\partial x} \\ \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + f u &= -g \frac{\partial \eta}{\partial y} \\ \frac{\partial \eta}{\partial t} + \frac{\partial}{\partial x}(hu) + \frac{\partial}{\partial y}(hv) &= 0\end{aligned}$$

Ekman pump (Northern Hemisphere)

$$\bar{w} = \frac{d}{2} \bar{\zeta}, \quad d = \left(\frac{2\nu_E}{f} \right)^{1/2} \quad \text{and} \quad w_{\text{Ek}} = \frac{1}{\rho_0 f} \left[\frac{\partial \tau^y}{\partial x} - \frac{\partial \tau^x}{\partial y} \right]$$

Equation sheet 2

- Shallow water equations for 2 layer model (linear, $f = f_0 + \beta_0 y$):

$$\begin{aligned} \frac{\partial u_1}{\partial t} - f v_1 &= -g \frac{\partial \eta}{\partial x} & \frac{\partial v_1}{\partial t} + f u_1 &= -g \frac{\partial \eta}{\partial y} \\ \frac{\partial u_2}{\partial t} - f v_2 &= -g \frac{\partial \eta}{\partial x} - g' \frac{\partial \eta}{\partial x} & \frac{\partial v_2}{\partial t} + f u_2 &= -g \frac{\partial \eta}{\partial y} - g' \frac{\partial \eta}{\partial y} \end{aligned}$$

$$\begin{aligned} \frac{\partial}{\partial t} (\eta - \alpha) + \frac{\partial}{\partial x} (H_1 u_1) + \frac{\partial}{\partial y} (H_1 v_1) &= 0 \\ \frac{\partial \eta}{\partial t} + \frac{\partial}{\partial x} [(H_2 - b) u_2] + \frac{\partial}{\partial y} [(H_2 - b) v_2] &= 0 \end{aligned}$$

Characteristic depth: $\bar{h} = \frac{H_1 H_2}{H_1 + H_2}$

- Generalised equation for barotropic planetary waves/topographic waves:

$$\frac{\partial \eta}{\partial t} - R^2 \frac{\partial}{\partial t} \nabla^2 \eta - \beta_0 R^2 \frac{\partial \eta}{\partial x} + \frac{\alpha_0 g}{f_0} \frac{\partial \eta}{\partial x} = 0$$

- Complex variable: $\phi \equiv \phi_r e^{i\theta} \equiv \phi_r + i\phi_i$; where $|\phi|^2 = \phi_r^2 + \phi_i^2$, $\tan \theta = \frac{\phi_i}{\phi_r}$

- QG Theory for continuously stratified fluid:

$$\frac{\partial \omega}{\partial z} = \frac{1}{\rho_0 f_0^2} \left[\frac{\partial}{\partial t} \nabla^2 p' + \frac{1}{\rho_0 f_0} J(p', \nabla^2 p') + \beta_0 \frac{\partial p'}{\partial x} \right]$$

$$\frac{\partial \rho'}{\partial t} + \frac{1}{\rho_0 f_0} J(p', \rho') - \frac{\rho_0 N^2}{g} \omega = 0$$

$$N^2 = -\frac{g}{\rho_0} \frac{d\bar{\rho}}{dz}$$

$$\frac{\partial q}{\partial t} + J(\psi, q) = 0, \quad q = \nabla^2 \psi + \frac{\partial}{\partial z} \left(\frac{f_0^2}{N^2} \frac{\partial \psi}{\partial z} \right) + \beta_0 y$$

- 2 layer QG model:

$$q_1 = \nabla^2 \psi_1 + \frac{1}{2R^2} (\psi_2 - \psi_1) + \beta_0 y$$

$$q_2 = \nabla^2 \psi_2 - \frac{1}{2R^2} (\psi_2 - \psi_1) + \beta_0 y$$

$$\omega = \frac{2f_0}{N^2 H} \left[\frac{\partial}{\partial t} (\psi_2 - \psi_1) + J(\psi_1, \psi_2) \right]$$

$$0 = \frac{f_0}{g} (\psi_2 - \psi_1)$$

$$\rho' = \frac{1}{g} \frac{(\psi_2 - \psi_1)}{H}$$

$$R = \frac{\sqrt{g'H}}{2f_0}$$

$$N^2 = \frac{2g'}{H}$$

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